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Working Group Summary: Pion-Nucleon Coupling Constant

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P.O.B. 9, FIN-00014 Helsinki, Finland***Abstract**

A brief introduction to different determinations of the πNN coupling constant is given, and some comments on the topics discussed in the working group are made.

INTRODUCTION

Since the birth of the Yukawa theory of the nuclear force in 1935 it was a challenge for the physics community to determine the coupling strength of the Yukawa meson to the nucleon. In 1947 the meson – pion – was finally discovered in cosmic ray emulsion experiments[1] and more systematic work to determine the πNN coupling constant could start. Conventionally[2] the pseudoscalar strength is denoted by g and the pseudovector coupling constant by f such that

$$f^2 = \left(\frac{M_\pi}{2m_p} \right)^2 \frac{g^2}{4\pi}, \quad (1)$$

where M_π is the charged pion mass and m_p is the proton mass. Other conventions concerning the nucleon mass and the factor 4π appear in the literature[3]. Reasonable estimates for the coupling strength were obtained even before the discovery of the pion and without detailed knowledge of the meson mass, e.g., Bethe was able to get an estimate $f^2 = 0.077 - 0.080$ already in 1940[4] on the basis of deuteron properties. The results of various determinations until 1980 are shown in Fig. 1. In the same figure very different techniques to determine f^2 are summarized. In the previous *MENU* symposium de Swart gave a review on the topic[3] and many of the references used in Fig. 1 can be found there. The values of

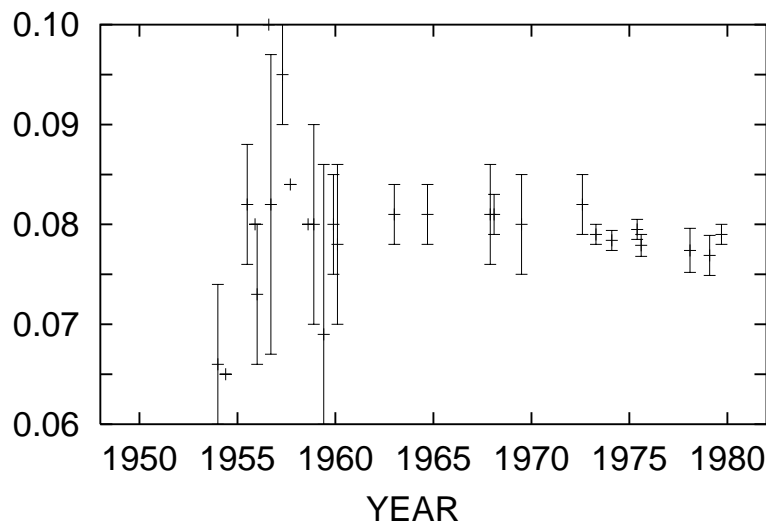
PION-NUCLEON COUPLING CONSTANT UNTIL 1980

Figure. 1. The values of the pion-nucleon coupling constant f^2 before 1980.

PION-NUCLEON COUPLING CONSTANT AFTER 1980

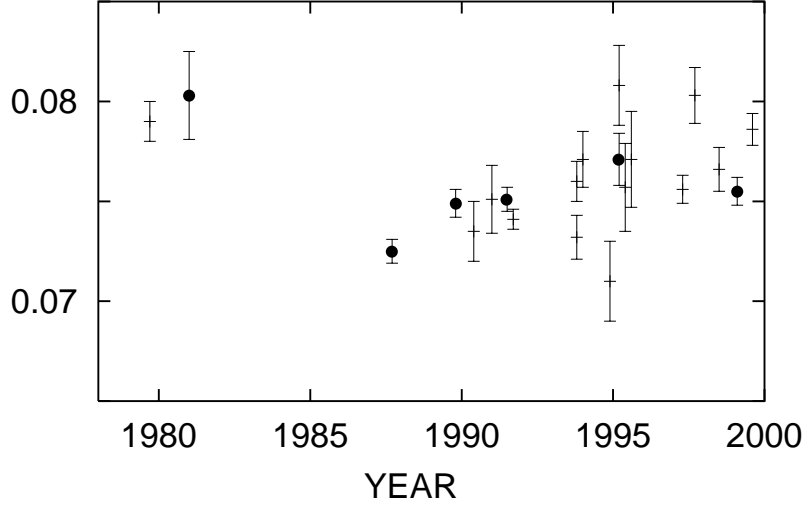


Figure. 2. The values of the pion-nucleon coupling constant f^2 after 1980 until the present. Neutral pion couplings are denoted by the solid dots, the remaining points refer to charged pion couplings or charge independent determinations.

f^2 stabilized for a long time[5–7] and only in the 90’s has the discussion of the value of the pion-nucleon coupling constant started again. In[5–7] fixed- t dispersion relations for πN were used. In the determinations displayed in Fig. 1 most of the data date back to the era before the meson factories, LAMPF, SIN and TRIUMF, which, in addition to performing experiments with pions, had programmes to study the NN interaction. In several analyses shown in Fig. 2 NN scattering data have been used to extract the πN coupling strength, i.e. one of the standard methods until the 60’s has been adopted again in more refined form. In this activity the Nijmegen group has played an important role[3]. Of course, in Fig. 2 many results from the meson factory πN experiments are included as well in the data bases used to determine f^2 .

The central issue in the discussion in the working group has been the scatter of the results of the determinations as shown in Fig. 2. The main questions involve the model dependence of different techniques, the effect of different pieces of data (partly conflicting), the error estimates, the electromagnetic corrections and other isospin violating effects. In the working group contributions were presented by Loiseau[8], Höhler[9] and Pavan[10], and brief commentaries by W.R. Gibbs, M. Birse and D.V. Bugg.

PROBLEMS IN EXTRACTING f^2

The particular issues raised in the discussion include:

- Electromagnetic corrections:
 - πN vs. NN; the different treatment of electromagnetic corrections for these two scattering processes gives a possibility to check the uncertainty due to these effects
 - corrections from Tromborg et al.[11] vs. Oades et al.[12]; the dispersion approach and potential model lead to differences which need checking
 - corrections at high energy; these need to be checked, Tromborg et al. calculated corrections only up to 655 MeV/c
- Lack of transparency of the analyses; the analyses contain large data bases and it is hard to clarify which pieces of information are the crucial ones in determining f^2

- The normalization of the np data is a problem in the (p, n) data analyses
- The determination of the s-wave isoscalar scattering length, a_{0+}^+ , from the π^-d level shift measurement suffers from some model dependence due to electromagnetic and absorption corrections
- Effective theory is not at present suitable for fixing the coupling constant. The problems relate mostly to the convergence of the chiral expansion or to the additional low-energy constants which are not known accurately enough. However, there might be a chance in a precise measurement of the induced pseudoscalar coupling constant, g_P , which would make an accurate determination of the pion-nucleon coupling constant possible[13]
- There is need for a new fixed- t analysis of πN scattering data which extends beyond the present limit of the VPI analysis, 2.1 GeV
- There is need for a new analysis of the forward dispersion relations of the NN system. The amount of NN data has increased considerably since the previous analysis thanks to the meson factories and SATURNE.

The GMO Sum Rule

The Goldberger-Miyazawa-Oehme sum rule (GMO)[14] provides a simple means to estimate the pion-nucleon coupling constant directly from measurable quantities, the πN isovector s-wave scattering length and total cross sections from the threshold to the highest energies. The method still has uncertainties, and will probably never be able to compete with other methods in precision, but the advantage is the possibility to relate the uncertainty in f^2 directly to the experimental errors.

The GMO sum rule is the result of the forward dispersion relation for the $D^-(=A^- + \nu B^-)$ amplitude taken at the physical threshold (the total laboratory energy $\omega = M_\pi$)

$$D^-(M_\pi) = \frac{8\pi f^2}{M_\pi[1 - (M_\pi^2/4m_p^2)]} + 4\pi M_\pi J^- = 4\pi(1+x)a_{0+}^-, \quad (2)$$

where

$$J^- = \frac{1}{2\pi^2} \int_0^\infty \frac{\sigma^-(k)}{\omega} dk \quad (3)$$

and $x = M_\pi/m_p$. The pion-nucleon coupling constant can now be extracted and the result is

$$\begin{aligned} f^2 &= \frac{1}{2} \left[1 - \left(\frac{x}{2} \right)^2 \right] [(1+x)M_\pi a_{0+}^- - M_\pi^2 J^-] \\ &= 0.5712(M_\pi a_{0+}^-) - 0.02488(J^-/\text{mb}). \end{aligned} \quad (4)$$

The isovector s-wave scattering length, a_{0+}^- , is accessible through experiment[19]. For the integral J^- several evaluations are displayed in Table 1. As can be seen from Fig. 3 there is potential sensitivity to details of the electromagnetic corrections especially around the Δ -resonance region near 0.3 GeV/c as well as to the treatment of the Δ^{++} , Δ^0 splitting. Making use of the isospin symmetry gives for the scattering length $a_{0+}^- = 0.0962 \pm 0.0071 M_\pi^{-1}$ [19] and taking Koch's value for J^- gives an estimate for the lower limit of the coupling constant f^2 with the result 0.0765. With more conservative errors for J^- the figure 0.0762 is obtained. With the remaining uncertainties in the treatment of various corrections this limit is not in real conflict with the results from other analyses.

Table 1. Values for the J^- integral.

Ref.	J^- (mb)
KH ('83)[2]	-1.058
Koch ('85)[15]	-1.077 ± 0.047
VPI ('92)[16]	-1.072
Gibbs ('98)[17]	-1.051
ELT ('99)[18]	-1.083 ± 0.025

ISOVECTOR COMBINATION OF TOTAL CROSS SECTIONS

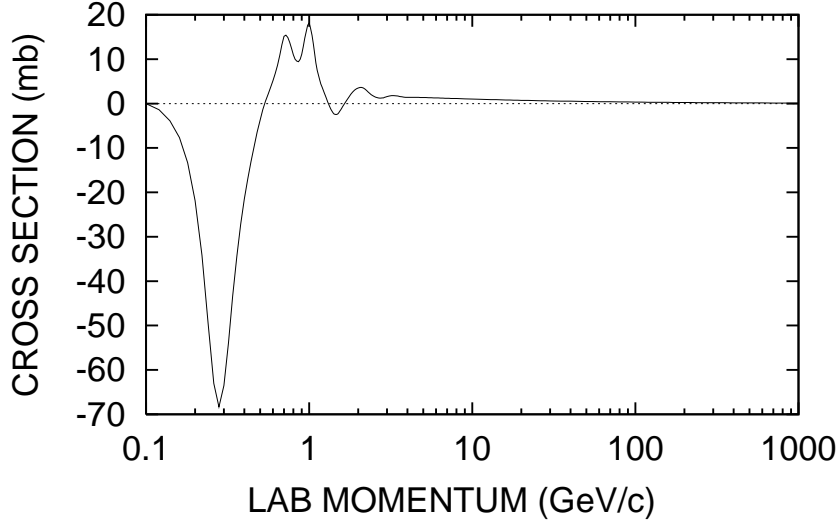


Figure. 3. The isovector combination, $\sigma^- = \frac{1}{2}(\sigma_{\pi^-p} - \sigma_{\pi^+p})$, of the π^-p and π^+p total cross sections[2]. Experimental data extend up to 350 GeV/c.

CONCLUDING REMARKS

The precision of the pion-nucleon experiments has now reached the level where a more careful treatment of the corrections, in particular of electromagnetic origin or due to the u - and d -quark mass difference, is necessary. These theoretical challenges have not yet been met, even though the theoretical tool, chiral perturbation theory, has now the capability to answer these questions. Work along these lines is in progress. In the lattice[20] and the QCD sum rule[21] frontiers the present accuracy for g is 20-30 % and it will take a while before this improves significantly.

Table 2 summarizes some recent values for f^2 displayed in Fig. 2. The table demonstrates the current trend, the favoured value for f^2 is slightly smaller than the standard one of Koch and Pietarinen[7]. However, there remains still quite a number of problems which need attention.

The Goldberger-Treiman discrepancy

$$\Delta_{\pi N} = 1 - \frac{m_p g_A}{F_\pi g}, \quad (5)$$

where F_π and g_A are the pion and neutron decay constants respectively, would be reduced from 4 % to 2 %, if f^2 changes from 0.079 to 0.076. In a recent SU(3) analysis[27] preference for a smaller Goldberger-Treiman discrepancy was found.

In the analysis of np scattering data at backward directions somewhat higher value for the coupling constant has been obtained[28], $f^2 = 0.0803 \pm 0.0014$. Discussion on

Table 2. Values for the pion-nucleon coupling constant f^2 from recent determinations.

Ref.	f^2	Method
KH ('80)[7]	0.079 ± 0.001	π N fixed- t
BM ('95)[22]	0.0757 ± 0.0022	NN data
Gibbs ('98)[17]	0.0756 ± 0.0007	GMO
Machner ('98)[23]	0.0760 ± 0.0011	symmetries
Matsinos ('98)[24]	0.0766 ± 0.0011	model fit
Nijmegen ('99)[25]	0.0756 ± 0.0004	pp PWA
ELT ('99)[18]	0.0786 ± 0.0008	GMO + π^-d
VPI ('99)[10,26]	0.0760 ± 0.0004	π N fixed- t

the problems in this field continues[29,30]. The spin transfer coefficients in pp scattering are also of interest, the preliminary indications are towards slightly smaller value for the coupling[31]. Machleidt has recently discussed[32] some additional problems with the deuteron properties and low-energy NN analyzing powers which indicate that no coherent picture is yet emerging.

ACKNOWLEDGEMENTS

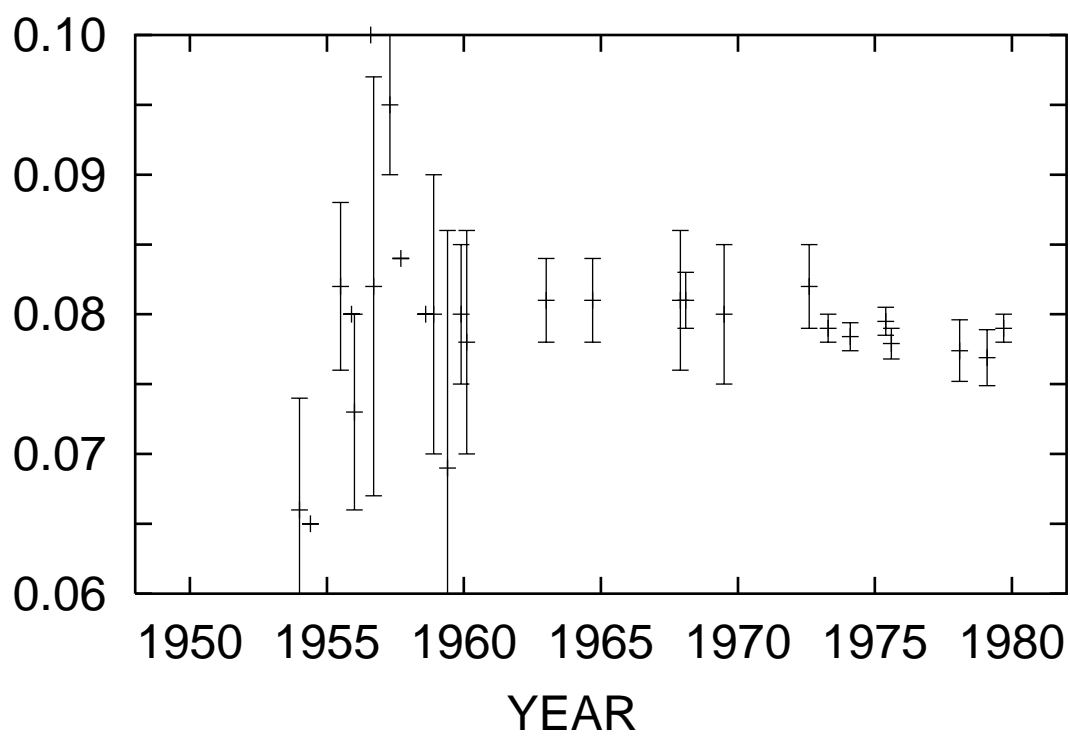
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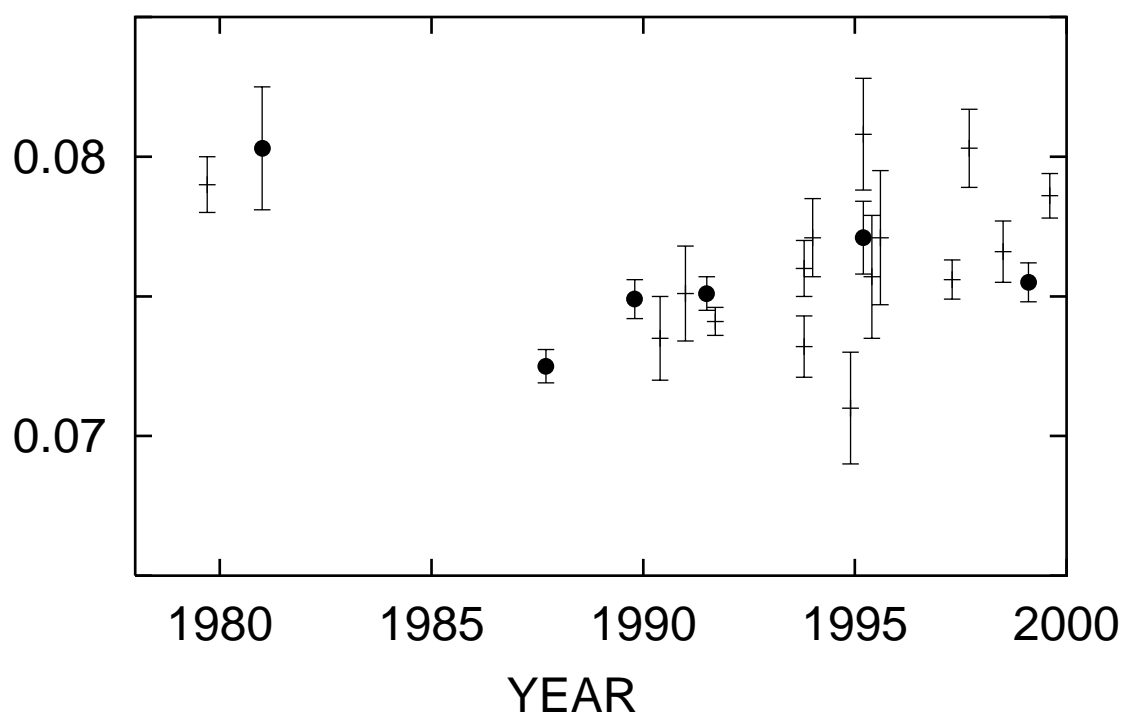
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